

# $^1\text{H}$ NMR Study of Ionic Motions in High Temperature Solid Phases of $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$

Hiroyuki Ishida

Department of Chemistry, Faculty of Science, Okayama University, Okayama 700-8530, Japan

Reprint requests to Prof. Dr. H. I.; E-mail: ishidah@cc.okayama-u.ac.jp

Z. Naturforsch. **55 a**, 412–414 (2000); received January 8, 2000

The reorientation of the tetrahedral complex anion  $\text{ZnCl}_4^{2-}$  and the self-diffusion of the cation in  $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$  were studied by  $^1\text{H}$  NMR spin-lattice relaxation time ( $^1\text{H}$   $T_1$ ) experiments. In the second highest-temperature phase, the temperature dependence of  $^1\text{H}$   $T_1$  observed at 8.5 MHz could be explained by a magnetic dipolar-electric quadrupolar cross relaxation between  $^1\text{H}$  and chlorine nuclei, and the activation energy of the anion motion was determined to be  $105\text{ kJ mol}^{-1}$ . In the highest-temperature phase, the activation energy of the self-diffusion of the cation was determined to be  $58\text{ kJ mol}^{-1}$  from the temperature and frequency dependence of  $^1\text{H}$   $T_1$ .

**Key words:** Nuclear Magnetic Resonance; Molecular Motion; Cross Relaxation;  $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$ .

Bis(methylammonium) tetrachlorozincate(II),  $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$ , forms at room temperature monoclinic crystals ( $\text{P2}_1/\text{a}$ ), containing discernible tetrahedral  $\text{ZnCl}_4^{2-}$  ions [1]. Pérez-Mato et al. reported the existence of two solid-solid phase transitions at 426 and 483 K from calorimetric, dielectric, thermal expansion, and optical measurements [2]. The transition at 426 K was also reported from Raman [3] and IR [4] spectra, but not found by DSC [5, 6], DTA, and  $^1\text{H}$  NMR [5] measurements. The transition at 483 K is accompanied by large enthalpy and entropy changes [1, 5, 6]. Previous  $^1\text{H}$  NMR studies revealed that the cation in the highest-temperature phase (Phase I) performs isotropic rotation and self-diffusion. The cation in the low temperature phases (Phase II and III) undergoes reorientation about its C-N bond axis [5]. However, the reorientational motion of the anion has not been clarified. In the present study, we measured the  $^1\text{H}$   $T_1$  at 8.5 MHz, which is approximately equal to an average value of  $^{35}\text{Cl}$  NQR resonance frequencies observed for eight salts containing tetrahedral  $\text{ZnCl}_4^{2-}$  ions [7], and showed that the motion of anions could be detected from the  $^1\text{H}$   $T_1$  through a cross relaxation mechanism between magnetic dipole ( $^1\text{H}$ ) and electric quadrupole ( $^{35}\text{Cl}$ ) energy levels. In addition, we measured  $^1\text{H}$   $T_1$  at two different resonance frequencies in Phase I to obtain more detailed information on the cationic dynamics.

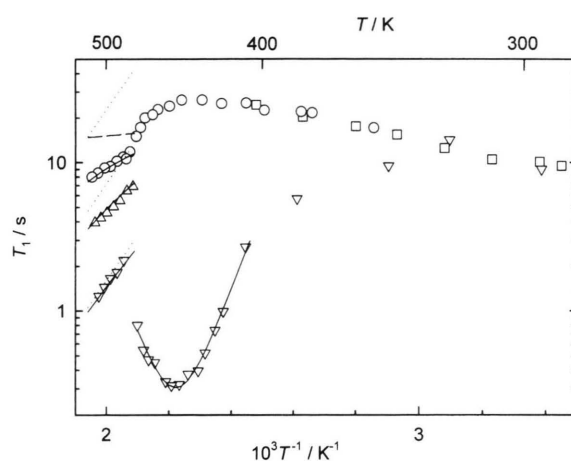


Fig. 1. Temperature dependence of  $^1\text{H}$   $T_1$  observed for  $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$  at 8.5 ( $\nabla$ ), 18 ( $\Delta$ ), 20 ( $\square$ ), and 32 ( $\circ$ ) MHz. Solid lines are best-fitted values calculated by (1) and (2). Dotted and dashed lines are calculated  $T_{1\text{DD}}$  and  $T_{1\text{SR}}$  values, respectively.

$(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$  was prepared and purified by the method reported in [5].  $^1\text{H}$  NMR spin-lattice relaxation times were measured at 18 and 8.5 MHz using the pulse spectrometer reported in [5].

The temperature dependence of  $T_1$  is shown in Fig. 1 together with the  $T_1$  data previously observed at 32 and 20 MHz [5]. Above 380 K in Phase II,  $T_1$  observed at 8.5 MHz decreased with increasing

0932–0784 / 00 / 0300–0412 \$ 06.00 © Verlag der Zeitschrift für Naturforschung, Tübingen · www.znaturforsch.com



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland Lizenz.

Zum 01.01.2015 ist eine Anpassung der Lizenzbedingungen (Entfall der Creative Commons Lizenzbedingung „Keine Bearbeitung“) beabsichtigt, um eine Nachnutzung auch im Rahmen zukünftiger wissenschaftlicher Nutzungsformen zu ermöglichen.

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

On 01.01.2015 it is planned to change the License Conditions (the removal of the Creative Commons License condition “no derivative works”). This is to allow reuse in the area of future scientific usage.

temperature and yielded a minimum of 0.31 s at ca. 450 K. This behavior is unexplainable by the magnetic dipole-dipole interactions between protons because the usual BPP theory using  $T_1$  data at 20 and 32 MHz predicts that the onset of decrease in  $T_1$  at 8.5 MHz is above 420 K, which is much higher than 380 K. Moreover, 8.5 MHz is approximately equal to an average value of  $^{35}\text{Cl}$  NQR resonance frequencies in tetrahedral  $\text{ZnCl}_4^{2-}$  ions, resulting in a strong cross relaxation between proton and chlorine nuclei. Therefore we believe that the relaxation mechanism at 8.5 MHz is mainly a modulation of  $^1\text{H}\dots^{35}\text{Cl}$  magnetic dipolar interactions due to anionic reorientations. In fact, Yamamoto et al. studied anionic motions in  $[(\text{CH}_3)_2\text{NH}_2]_2\text{ZnCl}_4$  by means of  $^{35}\text{Cl}$  NQR and found that above room temperature the  $^{35}\text{Cl}$  NQR spin-lattice relaxation time ( $T_{1Q}$ ) decreases rapidly with increasing temperature to less than 10 ms owing to the reorientation of the anion about its pseudo- $C_3$  axis [8]. Such a motion is expected to occur in  $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$  in the temperature region in question. When the resonance angular frequency of protons,  $\omega_H$ , is comparable to the quadrupole resonance frequency of chlorine nuclei,  $^1\text{H}$   $T_1$  is approximately given by [9]

$$T_1^{-1} = C\tau / (1 + (\omega_H - \langle\omega_{\text{Cl}}\rangle)^2\tau^2), \quad (1)$$

where  $\tau$ ,  $\omega_{\text{Cl}}$ , and  $\langle\ \rangle$  are the correlation time of anionic motion, the resonance angular frequency of chlorine nuclei, and the powder and isotope averages, respectively. Assuming an Arrhenius relationship for the activation energy  $E_a$  for the anionic motion:  $\tau = \tau_0 \exp(E_a/RT)$ , we fitted (1) to the  $T_1$  data, taking  $E_a$ ,  $C$ , and  $(\omega_H - \langle\omega_{\text{Cl}}\rangle)\tau_0$  as parameters. We obtained  $E_a$  of  $104 \pm 3$  kJ mol $^{-1}$ , which is comparable to that of 140 kJ mol $^{-1}$  obtained in  $[(\text{CH}_3)_2\text{NH}_2]_2\text{ZnCl}_4$  [7].

In Phase I,  $T_1$  decreased with increasing temperature. This decrease is attributable to the self-diffusion of the cation from the previous NMR study [5]. However, the gradient of the  $\log T_1$  vs.  $T^{-1}$  plots observed became gentler with increasing the Larmor

frequency, implying the presence of more than one relaxation mechanisms. A similar behavior of  $T_1$  was also reported for the CsCl-type ionic plastic phase of  $\text{CH}_3\text{NH}_3\text{X}$  ( $\text{X} = \text{NO}_3$  [10], I [11],  $\text{ClO}_4$  [12], Br [13]). According to the analysis of the  $T_1$  data of these salts, the present  $T_1$  values could also be expressed by the superposition of two components,  $T_{1\text{DD}}$  and  $T_{1\text{SR}}$ ,

$$T_1^{-1} = T_{1\text{DD}}^{-1} + T_{1\text{SR}}^{-1}, \quad (2)$$

where

$$T_{1\text{DD}}^{-1} = C_{\text{DD}}\omega_H^{-2}\tau_d^{-1}, \quad T_{1\text{SR}}^{-1} = C_{\text{SR}}\tau_r^{-1}. \quad (3)$$

$T_{1\text{DD}}$  denotes the relaxation time due to the magnetic dipole interaction among  $^1\text{H}$  nuclei modulated by the cationic self-diffusion.  $C_{\text{DD}}$  and  $\tau_d$  are the motional constant and the correlation time of cationic self-diffusion, respectively.  $T_{1\text{SR}}$  originates from the spin-rotation interaction due to the rapid uniaxial and/or isotropic rotation of the cation [14].  $C_{\text{SR}}$  and  $\tau_r$  are the motional constant and the correlation time of cationic rotation, respectively. In (2) and (3), we assume that  $\tau_r$  is much shorter than  $\tau_d$  and the conditions  $\omega_H\tau_d \gg 1$  and  $\omega_H\tau_r \ll 1$  are fulfilled. An Arrhenius-type temperature dependence can be approximated for the correlation time. Using (2) and (3),  $T_{1\text{DD}}$  and  $T_{1\text{SR}}$  was calculated separately, as shown in Figure 1. The activation energy evaluated for the cation self-diffusion was  $58 \pm 5$  kJ mol $^{-1}$ , which is much larger than 26 kJ mol $^{-1}$  previously estimated from  $T_1$  data measured only at 32 MHz [5]. On the other hand,  $E_a$  derived from  $T_{1\text{SR}}$  is  $4 \pm 2$  kJ mol $^{-1}$ . Thus, the spin-rotation interaction is due to the rapid uniaxial rotation of the cation, because this value is comparable to  $E_a$  of 3.8 kJ mol $^{-1}$  for the reorientation of  $\text{CH}_3\text{NH}_3^+$  about the C-N axis in  $(\text{CH}_3\text{NH}_3)_2\text{ZnCl}_4$  [5].

#### Acknowledgement

This work was supported by Grant-in-Aid for Scientific Research (B) (No. 10440208) from the Ministry of Education, Science, Sports and Culture, Japan.

- [1] B. Morosin and K. Emerson, *Acta Cryst.* **B32**, 294 (1976).
- [2] J. M. Pérez-Mato, J. L. Mañes, J. Fernández, J. Zúñiga, M. J. Tello, C. Socías, and M. A. Arriadiaga, *Phys. Stat. Sol. (a)* **68**, 29 (1981).
- [3] P. S. R. Prasad, *Phys. Stat. Sol. (a)* **149**, K13 (1995).
- [4] T. K. K. Srinivasan, M. Mylrajan, and J. B. Srinivasa Rao, *J. Raman Spectrosc.* **23**, 21 (1992).
- [5] H. Ishida, T. Iwachido, N. Hayama, R. Ikeda, M. Terashima, and D. Nakamura, *Z. Naturforsch.* **44a**, 741 (1989).
- [6] Y. Sakiyama, K. Horiuchi, and R. Ikeda, *J. Phys. Condens. Matter* **8**, 5345 (1996).
- [7] Nuclear Quadrupole Resonance Spectra (NQRS) Database (1999), Japan Association for International Chemical Information.
- [8] H. Yamamoto, A. Ishikawa, T. Asaji, and D. Nakamura, *Z. Naturforsch.* **45a**, 464 (1990).
- [9] Y. Furukawa, S. Gima, and D. Nakamura, *Ber. Bunsenges. Phys. Chem.* **89**, 863 (1985); Y. Tai, T. Asaji, R. Ikeda, and D. Nakamura, *Z. Naturforsch.* **44a**, 300 (1989), and references therein.
- [10] H. Ishida, R. Ikeda, and D. Nakamura, *J. Chem. Soc. Faraday Trans. 2*, **81**, 963 (1985).
- [11] H. Ishida, R. Ikeda, and D. Nakamura, *Bull. Chem. Soc. Japan* **59**, 915 (1986).
- [12] H. Ishida, R. Ikeda, and D. Nakamura, *Bull. Chem. Soc. Japan* **60**, 467 (1987).
- [13] M. Tansho, D. Nakamura, and R. Ikeda, *Ber. Bunsenges. Phys. Chem.* **95**, 1643 (1991).
- [14] P. S. Hubbard, *Phys. Rev.* **131**, 1155 (1963).